

Roadmap for extreme-scale simulations

F.Xavier Trias¹, Àdel Alsalti-Baldellou^{1,2}, Assensi Oliva¹

¹Heat and Mass Transfer Technological Center, Technical University of Catalonia ²Department of Information Engineering, University of Padova, Italy





Roadmap for extreme-scale simulations: on the evolution of Poisson solvers

F.Xavier Trias¹, Àdel Alsalti-Baldellou^{1,2}, Assensi Oliva¹

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Motivation 00	Two competing effects	Residual of Poisson's equation	Solver convergence	Results 00	Conclusions 00

Contents



- 2 Two competing effects
- 3 Residual of Poisson's equation
- 4 Solver convergence

5 Results



10 EFlop/s Exascale 1 EFlop/s 100 PFlop/s 10 PFlop/s And in case of the Petascale 1 PFlop/s Performance a start and the 100 TFlop/s 10 TFlop/s Terascale 1 TFlop/s 100 GFlop/s 10 GFlop/s ¢* ---- Sum **---** #1 - #500 1 GFlop/s 100 MFlop/s 1990 1995 2000 2005 2010 2015 2020 2025

Projected Performance Development

Source: www.top500.org













			~10 y	ears ~5 y	ears	
	PetaFLOPS		#1 in LINPACK	#1 in HPCG	Cutting-edge CFD simulation	'Routine' CFD simulation
Zetta	106		2037	2047	2052	
Exa	10 ³	years	2022 (Frontier)	2032	2037	
Peta	1	ars 14	2008 (Roadrunner)	2018 (Summit)	2023	
Tera	10-3	11 ye	1997 (ASCI Red)	No data		





Motivation ●0	Two competing effects	Residual of Poisson's equation	Solver convergence	Results 00	Conclusions

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					Cutting-edge	'Routine' CED
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Research question :

• Will the **complexity** of numerically solving **Poisson's equation** increase or decrease for **very large scale DNS/LES** simulations of incompressible turbulent flows?



DNS¹ of air-filled Rayleigh–Bénard convection at $Ra = 10^8$ and 10^{10}

¹B.Sanderse, F.X.Trias. *Energy-consistent discretization of viscous dissipation with application to natural convection flow*. (https://arxiv.org/abs/2307.10874)











Step 1:
$$\frac{\vec{u}^{p} - \vec{u}^{n}}{\Delta t} = \frac{3}{2} \vec{R} (\vec{u}^{n}) - \frac{1}{2} \vec{R} (\vec{u}^{n-1})$$
Step 2:
$$\nabla^{2} p^{n+1} = \frac{1}{\Delta t} \nabla \cdot \vec{u}^{p}$$

$$\int_{At} \frac{3}{2} \vec{R} (\vec{u}^{n}) - \frac{1}{2} \vec{R} (\vec{u}^{n-1}) - \nabla p^{n+1}$$
Semi-discrete (just in time)
NS equations
$$\int_{NS} \frac{\vec{u}^{n+1} - \vec{u}^{n}}{\Delta t} = \frac{3}{2} \vec{R} (\vec{u}^{n}) - \frac{1}{2} \vec{R} (\vec{u}^{n-1}) - \nabla p^{n+1}$$

$$\nabla \cdot \vec{u}^{n+1} = 0$$



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$$\frac{\vec{u}^{p} - \vec{u}^{n}}{\Delta t} = \frac{3}{2} \vec{R} (\vec{u}^{n}) - \frac{1}{2} \vec{R} (\vec{u}^{n-1})$$
Step 2:
$$\nabla^{2} p^{n+1} = \frac{1}{\Delta t} \nabla \cdot \vec{u}^{p}$$
Step 3:
$$\vec{u}^{n+1} = \vec{u}^{p} - \Delta t \nabla p^{n+1}$$

$$\int_{u}^{\Delta t} \vec{u}^{n+1} = \frac{3}{2} \vec{R} (\vec{u}^{n}) - \frac{1}{2} \vec{R} (\vec{u}^{n-1}) - \nabla p^{n+1}$$
Semi-discrete (just in time)
NS equations
$$\int_{u}^{u} \vec{u}^{n+1} = \vec{u}^{n+1} - \vec{u}^{n} = \frac{3}{2} \vec{R} (\vec{u}^{n}) - \frac{1}{2} \vec{R} (\vec{u}^{n-1}) - \nabla p^{n+1}$$

$$\nabla \cdot \vec{u}^{n+1} = 0$$

Research question:

 Will the complexity of numerically solving Poisson's equation increase or decrease for very large scale DNS/LES simulations of incompressible turbulent flows?

$$\left(\nabla^2 p^{n+1} = \frac{1}{\Delta t} \nabla \cdot \vec{u}^p\right)$$

Two competing effects: who (if any) will eventually win?

Re
$$\uparrow$$
 $\Delta x \checkmark \longrightarrow N_x \uparrow \longrightarrow$ Larger system \checkmark
 $\Delta t \checkmark \longrightarrow$ Better initial guess \uparrow



Research question:

 Will the complexity of numerically solving Poisson's equation increase or decrease for very large scale DNS/LES simulations of incompressible turbulent flows?

 $Ra = 10^{8}$



²F.Dabbagh, F.X.Trias, A.Gorobets, A.Oliva. *Flow topology dynamics in a 3D phase space for turbulent Rayleigh-Bénard convection*, **Phys.Rev.Fluids**, 5:024603, 2020.



Research question:

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 Will the complexity of numerically solving Poisson's equation increase or decrease for very large scale DNS/LES simulations of incompressible turbulent flows?



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Smaller and smaller, but how much?









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Two competing effects: who (if any) will eventually win?

Re
$$\uparrow$$
 $\Delta x \downarrow \longrightarrow N_x \uparrow \longrightarrow$ Larger system \downarrow
 $\Delta t \downarrow \longrightarrow$ Better initial guess \uparrow

In summary:

$$\frac{1}{N_x^{K41}} = \frac{\Delta x}{L_x} \sim \frac{\eta}{l} \propto \text{Re}^{-3/4}$$

$$\alpha = -1/2 \quad (\text{ K41 or diffusion dominated })$$

$$\frac{\Delta t}{t_l} \sim \text{Re}^{\alpha} \qquad \alpha = -3/4 \quad (\text{ convection dominated })$$

Motivation 00	Two competing effects	Residual of Poisson's equation ●00	Solver convergence	Results 00	Conclusions

$$\nabla^{2} p^{n+1} = \frac{1}{\Delta t} \nabla \cdot \vec{u}^{p}$$

$$\downarrow \text{Initial guess} \Rightarrow p^{n}$$

$$r^{o} = \nabla^{2} p^{n} - \frac{1}{\Delta t} \nabla \cdot u^{p,n+1}$$

Motivation Two competing effects Residual of Poisson's equation Solver convergence Results Convergence 00 000 <	nclusions
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$$\nabla^{2} p^{n+1} = \frac{1}{\Delta t} \nabla \cdot \vec{u}^{p}$$
Initial guess $\rightarrow p^{n}$

$$\downarrow$$

$$r^{o} = \nabla^{2} p^{n} - \frac{1}{\Delta t} \nabla \cdot u^{p,n+1} = \frac{1}{\Delta t} \nabla \cdot u^{p,n} - \frac{1}{\Delta t} \nabla \cdot u^{p,n+1} \approx \frac{\partial \nabla \cdot u^{p}}{\partial t} = \nabla \cdot \frac{\partial u^{p}}{\partial t}$$

Motivation 00	Two competing effects	Residual of Poisson's equation ●00	Solver convergence	Results 00	Conclusions

$$\nabla^{2} p^{n+1} = \frac{1}{\Delta t} \nabla \cdot \vec{u}^{p}$$
Initial guess $\Rightarrow p^{n}$

$$r^{o} = \nabla^{2} p^{n} - \frac{1}{\Delta t} \nabla \cdot u^{p,n+1} = \frac{1}{\Delta t} \nabla \cdot u^{p,n} - \frac{1}{\Delta t} \nabla \cdot u^{p,n+1} \approx \frac{\partial \nabla \cdot u^{p}}{\partial t} = \nabla \cdot \frac{\partial u^{p}}{\partial t}$$

$$\tilde{r}^{o} = \nabla^{2} \tilde{p}^{n} - \nabla \cdot u^{p,n+1} \approx \nabla \cdot u^{p,n} - \nabla \cdot u^{p,n+1} \approx \Delta t \frac{\partial \nabla \cdot u^{p}}{\partial t} = \Delta t \nabla \cdot \frac{\partial u^{p}}{\partial t}$$
Initial guess $\Rightarrow \tilde{p}^{n} = \Delta t p^{n}$

$$\nabla^{2} \tilde{p}^{n+1} = \nabla \cdot \vec{u}^{p}$$

Motivation 00	Two competing effects	Residual of Poisson's equation ○●○	Solver convergence	Results 00	Conclusions

$$\nabla^{2} p^{n+1} = \frac{1}{\Delta t} \nabla \cdot \vec{u}^{p}$$
Initial guess $\Rightarrow p^{n}$

$$r^{o} \approx \frac{\partial \nabla \cdot u^{p}}{\partial t}$$

$$\tilde{r}^{o} \approx \Delta t \frac{\partial \nabla \cdot u^{p}}{\partial t}$$
Initial guess $\Rightarrow \tilde{p}^{n} = \Delta t p^{n}$

$$\nabla^{2} \tilde{p}^{n+1} = \nabla \cdot \vec{u}^{p}$$



$$\nabla^{2} p^{n+1} = \frac{1}{\Delta t} \nabla \cdot \vec{u}^{p}$$
Initial guess $\Rightarrow p^{n}$

$$Q_{G} - criterion$$

$$P^{o} \approx \frac{\partial \nabla \cdot u^{p}}{\partial t}$$
What is $\nabla \cdot u^{p}$?
$$\nabla \cdot u^{p} = \sum (n - \Delta t \nabla \cdot (u^{n} \cdot \nabla u^{n}) + v \Delta \nabla \cdot \vec{u}^{n} = 2\Delta t Q_{G}$$

$$P^{o} \approx \Delta t \frac{\partial \nabla \cdot u^{p}}{\partial t}$$
Initial guess $\Rightarrow p^{n} = \Delta t p^{n}$

$$\nabla^{2} \tilde{p}^{n+1} = \nabla \cdot \vec{u}^{p}$$

Motivation 00	Two competing effects	Residual of Poisson's equation ○●○	Solver convergence	Results 00	Conclusions

$$\nabla^{2} p^{n+1} = \frac{1}{\Delta t} \nabla \cdot \vec{u}^{p}$$
Initial guess $\Rightarrow p^{n}$

$$Q_{G} - criterion$$

$$r^{o} \approx \frac{\partial \nabla \cdot u^{p}}{\partial t} = 2 \Delta t \frac{\partial Q_{G}}{\partial t}$$

$$R_{G} = det(G) = \frac{1}{3} tr(G^{3})$$

$$\overline{r}^{o} \approx \Delta t \frac{\partial \nabla \cdot u^{p}}{\partial t} = 2 \Delta t^{2} \frac{\partial Q_{G}}{\partial t}$$

$$Q_{G} = -\frac{1}{2} tr(G^{2}) \text{ where } G = \nabla u^{n}$$

Exact equations for restricted Euler:

$$\frac{d\mathbf{Q}_G}{dt} = -3\mathbf{R}_G \longrightarrow \frac{\partial \mathbf{Q}_G}{\partial t} = -(\mathbf{u} \cdot \nabla)\mathbf{Q}_G - 3\mathbf{R}_G$$



$$\nabla^{2} p^{n+1} = \frac{1}{\Delta t} \nabla \cdot \vec{u}^{p}$$
Initial guess $\Rightarrow p^{n}$

$$Q_{G} - criterion$$

$$r^{o} \approx \frac{\partial \nabla \cdot u^{p}}{\partial t} = 2 \Delta t \frac{\partial Q_{G}}{\partial t} \approx -2 \Delta t \{(u \cdot \nabla) Q_{G} + 3 R_{G}\}$$

$$\tilde{r}^{o} \approx \Delta t \frac{\partial \nabla \cdot u^{p}}{\partial t} = 2 \Delta t^{2} \frac{\partial Q_{G}}{\partial t} \approx -2 \Delta t^{2} \{(u \cdot \nabla) Q_{G} + 3 R_{G}\}$$

Exact equations for restricted Euler:

$$\frac{dQ_G}{dt} = -3R_G \longrightarrow \frac{\partial Q_G}{\partial t} = -(u \cdot \nabla)Q_G - 3R_G$$

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	vergence Results Conclusions
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$$\nabla^{2} p^{n+1} = \frac{1}{\Delta t} \nabla \cdot \vec{u}^{p}$$
Initial guess $\Rightarrow p^{n}$

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Initial guess $\Rightarrow \tilde{p}^{n} = \Delta t p^{n}$

$$\nabla^{2} \tilde{p}^{n+1} = \nabla \cdot \vec{u}^{p}$$

Motivation 00	Two competing effects	Residual of Poisson's equation 00●	Solver convergence	Results 00	Conclusions 00
Residu	al of Poisson	's equation in Fo	ourier space		
In r	summary: ${}^{o} \approx \frac{\partial \nabla \cdot u^{p}}{\partial t} = 2\Delta$	$t^{p}\frac{\partial Q_{G}}{\partial t}\approx -2\Delta t^{p}\{(u)\}$	$(\cdot \nabla) \mathbf{Q}_{\mathbf{G}} + 3 \mathbf{R}_{\mathbf{G}}$	$\left\{ \nabla^{2} p^{n+1} = \frac{1}{2} \right\} p = \frac{1}{2}$	$\frac{\frac{1}{\Delta t} \nabla \cdot \vec{u}^{p}}{= \{1, 2\}}$ $= \nabla \cdot \vec{u}^{p}$

$$\frac{\Delta t}{t_{l}} \sim \operatorname{Re}^{\alpha} \begin{cases} \alpha = -1/2 \text{ (K41 or diffusion dominated)} \\ \alpha = -3/4 \text{ (convection dominated)} \end{cases}$$

$$\frac{1}{N_x^{K41}} = \frac{\Delta x}{L_x} \sim \frac{\eta}{l} \propto \text{Re}^{-3/4}$$

Residual of Poisson's equation in Fourier space

In summary:

$$r^{o} \approx \frac{\partial \nabla \cdot u^{p}}{\partial t} = 2 \Delta t^{p} \frac{\partial Q_{G}}{\partial t} \approx -2 \Delta t^{p} \{ (u \cdot \nabla) Q_{G} + 3 R_{G} \}$$

$$p = \{1, 2\}$$

$$p = \{1, 2\}$$

$$\nabla^{2} \tilde{p}^{au} = \nabla \cdot \tilde{u}^{p}$$
Hypothesis:
(inertial range)

$$\left(\frac{\partial \hat{Q}_{G}}{\partial t} \right)_{k} \propto k^{\beta} \longrightarrow \hat{r}_{k}^{o} \propto \Delta t^{p} k^{\beta} \sim \operatorname{Re}^{p\alpha} k^{\beta} = \operatorname{Re}^{\tilde{\alpha}} k^{\beta}$$

$$\frac{\Delta t}{t_{l}} \sim \operatorname{Re}^{\alpha} \left\{ \begin{array}{c} \alpha = -1/2 & (\text{ K41 or diffusion dominated }) \\ \alpha = -3/4 & (\text{ convection dominated }) \end{array} \right\}$$

$$\frac{1}{N_{x}^{K41}} = \frac{\Delta x}{L_{x}} \sim \frac{\eta}{l} \propto \operatorname{Re}^{-3/4}$$

Residual of Poisson's equation in Fourier space

In summary:

$$r^{o} \approx \frac{\partial \nabla \cdot u^{p}}{\partial t} = 2\Delta t^{p} \frac{\partial Q_{G}}{\partial t} \approx -2\Delta t^{p} \{(u \cdot \nabla) Q_{G} + 3R_{G}\} \qquad p = \{1, 2\}$$

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Hypothesis:
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$$\frac{\Delta t}{t_{l}} \sim \operatorname{Re}^{\alpha} \left\{ \begin{array}{c} \alpha = -1/2 & (\text{ K41 or diffusion dominated }) \\ \alpha = -3/4 & (\text{ convection dominated }) \end{array} \right\}$$
Parseval's theorem

$$\left\| r \right\|^{2} = \int_{\Omega} r^{2} dV = \int_{1}^{k_{max}} \hat{r}_{k}^{2} dk$$

Residual of Poisson's equation in Fourier space



Motivation 00	Two competing effects	Residual of Poisson's equation	Solver convergence ●0	Results 00	Conclusions

$$\|\boldsymbol{r}^{n}\|^{2} = \int_{1}^{k_{max}} (\hat{\omega}_{k}^{n} \hat{\boldsymbol{r}}_{k}^{0})^{2} dk \approx \int_{1}^{\operatorname{Re}^{34}} \hat{\omega}_{k}^{2n} \operatorname{Re}^{2\widetilde{\alpha}} k^{2\beta} dk$$
$$\hat{\omega} = \frac{\hat{\boldsymbol{r}}_{k}^{n+1}}{\hat{\boldsymbol{r}}_{k}^{n}} \int_{1}^{\infty} (\hat{\boldsymbol{r}}_{k}^{0} \propto \Delta t^{p} k^{\beta} \sim \operatorname{Re}^{p\alpha} k^{\beta} = \operatorname{Re}^{\widetilde{\alpha}} k^{\beta})^{2} dk$$

$$||\mathbf{r}||^{2} = \int_{\Omega} r^{2} dV = \int_{1}^{k_{max}} \hat{\mathbf{r}}_{k}^{2} dk \approx \int_{1}^{\text{Re}^{3/4}} \hat{\mathbf{r}}_{k}^{2} dk$$

Motivation 00	Two competing effects	Residual of Poisson's equation	Solver convergence ●0	Results 00	Conclusions

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Motivation 00	Two competing effects	Residual of Poisson's equation	Solver convergence ●0	Results 00	Conclusions

Motivation 00	Two competing effects	Residual of Poisson's equation	Solver convergence ●0	Results 00	Conclusions

Motivation 00	Two competing effects	Residual of Poisson's equation	Solver convergence ○●	Results 00	Conclusions



Motivation Two competing effects Residual of Poisson's equation	Solver convergence ○●	Results 00	Conclusions
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Motivation 00	Two competing effects	Residual of Poisson's equation	Solver convergence ○●	Results 00	Conclusions



Motivation 00	Two competing effects	Residual of Poisson's equation	Solver convergence ○●	Results 00	Conclusions





10

Re₂ ≈ 433 (1024³) from https://turbulence.pha.jhu.edu/

100

10⁻⁶

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Motivation 00	Two competing effects	Residual of Poisson's equation	Solver convergence	Results ●0	Conclusions

Kolmogorov theory predictions



Kolmogorov theory predictions















Motivation 00	Two competing effects	Residual of Poisson's equation	Solver convergence	Results 00	Conclusions ●0
Conclu	ding remarks				

• **Two competing effects** on the convergence of Poisson's equation have been identified.

Motivation 00	Two competing effects	Residual of Poisson's equation	Solver convergence	Results 00	Conclusions ●○

Concluding remarks

- **Two competing effects** on the convergence of Poisson's equation have been identified.
- The {α, β} phase space is divided in two regions depending on the solver convergence.



Concluding remarks

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- The {α, β} phase space is divided in two regions depending on the solver convergence.
- Numerical results match well with the developed theory prediction $\beta \approx 11/6$





Concluding remarks

- Two competing effects on the convergence of Poisson's equation have been identified.
- The {α, β} phase space is divided in two regions depending on the solver convergence.
- Numerical results match well with the developed theory prediction $\beta \approx 11/6$



- Extend the analysis up to $Re_\lambda \approx 1200$ using 8192^3 nodes
- Analysis of more complex flows







Thank you for your attendance

