



Centre Tecnològic de Transferència de Calor
UNIVERSITAT POLITÈCNICA DE CATALUNYA



Exa, zetta, yotta and beyond

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Exa, zetta, yotta and beyond: on the evolution of Poisson solvers

F.Xavier Trias¹, Àdel Alsalti-Baldellou^{1,2}, Assensi Oliva¹

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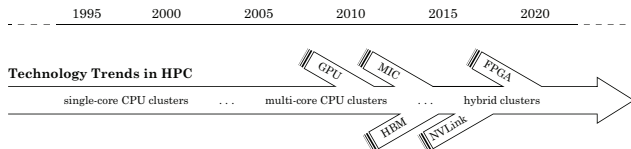
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- 3 Residual of Poisson's equation
- 4 Solver convergence
- 5 Results
- 6 Conclusions

Motivation

Research question #1:

- How can we develop **portable** and **efficient** CFD codes for large-scale simulations on modern supercomputers?



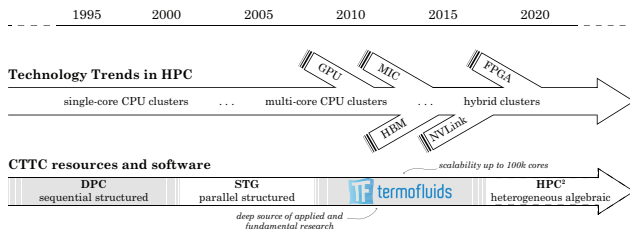
¹X.Álvarez, A.Gorobets, F.X.Trias. *A hierarchical parallel implementation for heterogeneous computing. Application to algebra-based CFD simulations on hybrid supercomputers.* **Computers & Fluids**, 214:104768, 2021.

²Á.Alsalti-Baldellou, X.Álvarez-Farré, F.X.Trias, A.Oliva. *Exploiting spatial symmetries for solving Poisson's equation.* **Journal of Computational Physics**, 486:112133, 2023.

Motivation

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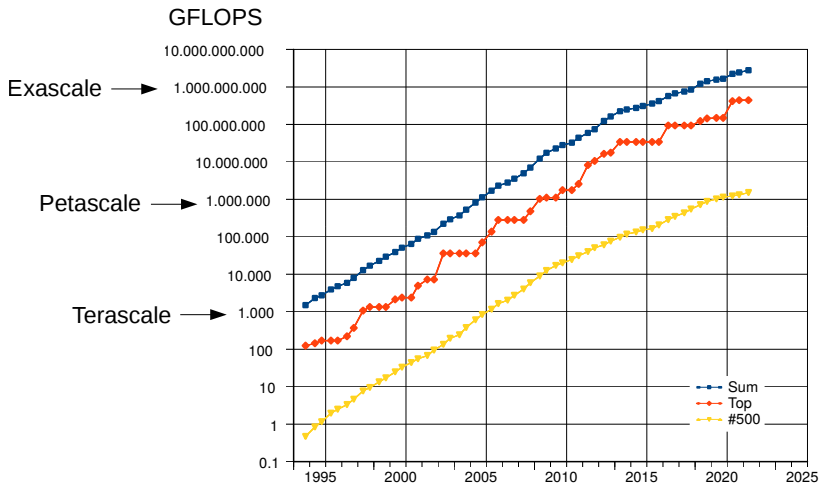


HPC²: portable, algebra-based framework for heterogeneous computing is being developed¹. Traditional stencil-based data and sweeps are replaced by algebraic structures (sparse matrices and vectors) and kernels. SpMM-based strategies to increase the arithmetic intensity are being considered².

¹X.Álvarez, A.Gorobets, F.X.Trias. *A hierarchical parallel implementation for heterogeneous computing. Application to algebra-based CFD simulations on hybrid supercomputers.* **Computers & Fluids**, 214:104768, 2021.

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Tera, Peta, Exa,..., Zetta, Yotta

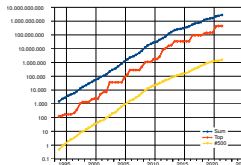


Tera, Peta, Exa, ..., Zetta, Yotta

~10 years



	PetaFLOPS		#1 in LINPACK	#1 in HPCG	Cutting-edge CFD simulation	'Routine' CFD simulation
Zetta	10^6					
Exa	10^3	14 years ↑	2022 (Frontier)			
Peta	1	11 years ↑	2008 (Roadrunner)	2018 (Summit)		
Tera	10^3		1997 (ASCI Red)	No data		

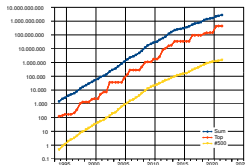


Tera, Peta, Exa, ..., Zetta, Yotta

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	PetaFLOPS		#1 in LINPACK	#1 in HPCG	Cutting-edge CFD simulation	'Routine' CFD simulation
Zetta	10^6		2037	2047		
Exa	10^3	14 years ↑	2022 (Frontier)	2032		
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Tera	10^3		1997 (ASCI Red)	No data		



Tera, Peta, Exa, ..., Zetta, Yotta

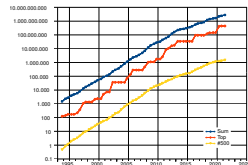
~10 years



~5 years



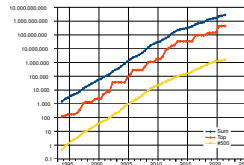
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Tera, Peta, Exa, ..., Zetta, Yotta

	PetaFLOPS		#1 in LINPACK	#1 in HPCG	Cutting-edge CFD simulation	'Routine' CFD simulation
Zetta	10^6		2037	2047	2052	2062
Exa	10^3	14 years ↑	2022 (Frontier)	2032	2037	2047
Peta	1	11 years ↑	2008 (Roadrunner)	2018 (Summit)	2023	2033
Tera	10^{-3}		1997 (ASCI Red)	No data		

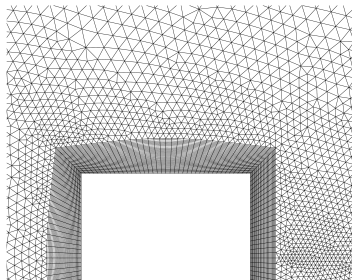
~10 years → ~5 years → ~10 years →



Motivation

Research question #2:

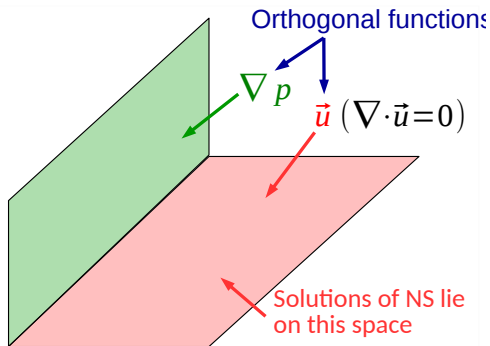
- Will the **complexity** of numerically solving **Poisson's equation** **increase** or **decrease** for **very large scale DNS/LES** simulations of incompressible turbulent flows?



DNS³ of the turbulent flow around a square cylinder at $Re = 22000$

³F.X.Trias, A.Gorobets, A.Oliva. *Turbulent flow around a square cylinder at Reynolds number 22000: a DNS study*, **Computers&Fluids**, 123:87-98, 2015.

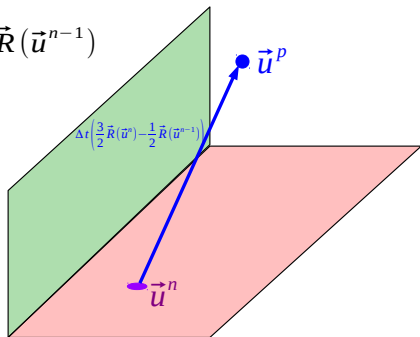
Poisson's equation: a quick reminder



$$\left. \begin{array}{l} \text{Semi-discrete} \\ \text{(just in time)} \\ \text{NS equations} \end{array} \right\} \begin{cases} \frac{\vec{u}^{n+1} - \vec{u}^n}{\Delta t} = \frac{3}{2} \vec{R}(\vec{u}^n) - \frac{1}{2} \vec{R}(\vec{u}^{n-1}) - \nabla p^{n+1} \\ \nabla \cdot \vec{u}^{n+1} = 0 \end{cases}$$

Poisson's equation: a quick reminder

$$\text{Step 1: } \frac{\vec{u}^p - \vec{u}^n}{\Delta t} = \frac{3}{2} \vec{R}(\vec{u}^n) - \frac{1}{2} \vec{R}(\vec{u}^{n-1})$$

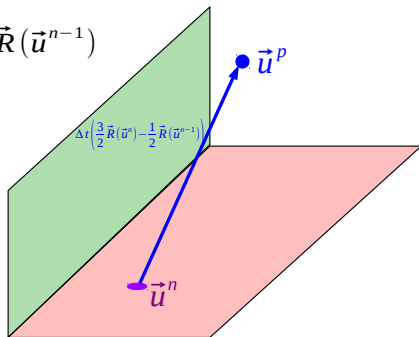


$$\text{Semi-discrete (just in time) NS equations} \left\{ \begin{array}{l} \frac{\vec{u}^{n+1} - \vec{u}^n}{\Delta t} = \frac{3}{2} \vec{R}(\vec{u}^n) - \frac{1}{2} \vec{R}(\vec{u}^{n-1}) - \nabla p^{n+1} \\ \nabla \cdot \vec{u}^{n+1} = 0 \end{array} \right.$$

Poisson's equation: a quick reminder

Step 1:
$$\frac{\vec{u}^p - \vec{u}^n}{\Delta t} = \frac{3}{2} \vec{R}(\vec{u}^n) - \frac{1}{2} \vec{R}(\vec{u}^{n-1})$$

Step 2:
$$\nabla^2 p^{n+1} = \frac{1}{\Delta t} \nabla \cdot \vec{u}^p$$



Semi-discrete (just in time) NS equations

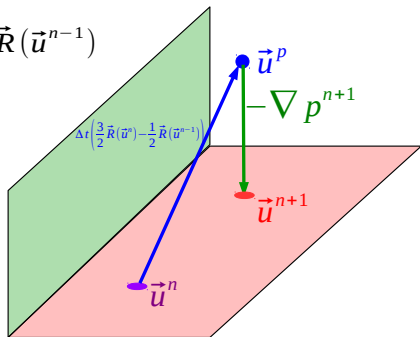
$$\left\{ \begin{array}{l} \frac{\vec{u}^{n+1} - \vec{u}^n}{\Delta t} = \frac{3}{2} \vec{R}(\vec{u}^n) - \frac{1}{2} \vec{R}(\vec{u}^{n-1}) - \nabla p^{n+1} \\ \nabla \cdot \vec{u}^{n+1} = 0 \end{array} \right.$$

Poisson's equation: a quick reminder

Step 1: $\frac{\vec{u}^p - \vec{u}^n}{\Delta t} = \frac{3}{2} \vec{R}(\vec{u}^n) - \frac{1}{2} \vec{R}(\vec{u}^{n-1})$

Step 2: $\nabla^2 p^{n+1} = \frac{1}{\Delta t} \nabla \cdot \vec{u}^p$

Step 3: $\vec{u}^{n+1} = \vec{u}^p - \Delta t \nabla p^{n+1}$



Semi-discrete (just in time) NS equations $\left\{ \begin{array}{l} \frac{\vec{u}^{n+1} - \vec{u}^n}{\Delta t} = \frac{3}{2} \vec{R}(\vec{u}^n) - \frac{1}{2} \vec{R}(\vec{u}^{n-1}) - \nabla p^{n+1} \\ \nabla \cdot \vec{u}^{n+1} = 0 \end{array} \right.$

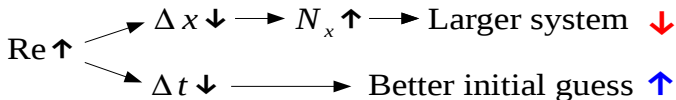
Poisson's equation: getting more tough or not?

Research question #2:

- Will the **complexity** of numerically solving **Poisson's equation** **increase** or **decrease** for **very large scale DNS/LES** simulations of incompressible turbulent flows?

$$\nabla^2 p^{n+1} = \frac{1}{\Delta t} \nabla \cdot \vec{u}^p$$

Two competing effects: who (if any) will eventually win?

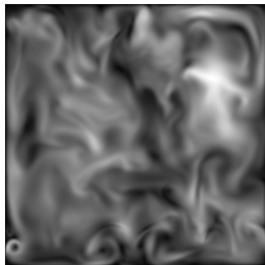


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$$Ra = 10^8$$

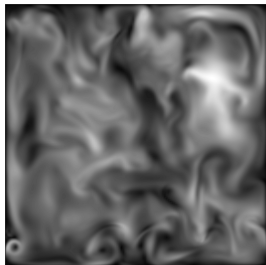


Poisson's equation: getting more tough or not?

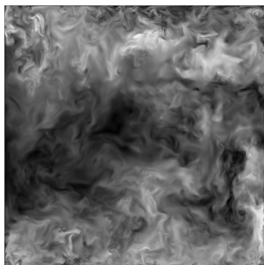
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$Ra = 10^8$



$Ra = 10^{10}$

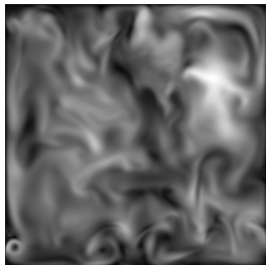


Poisson's equation: getting more tough or not?

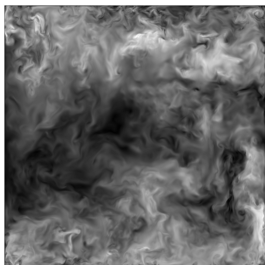
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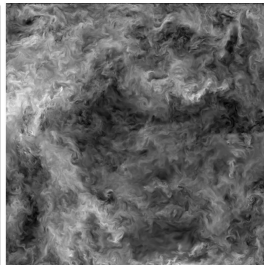
$Ra = 10^8$



$Ra = 10^{10}$



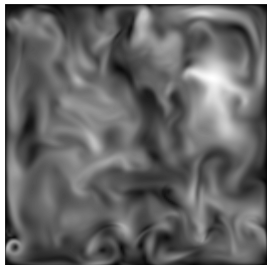
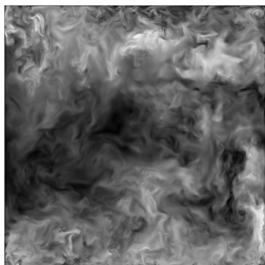
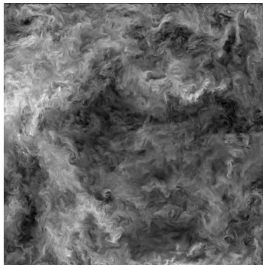
$Ra = 10^{11}$



Poisson's equation: getting more tough or not?

Research question #2:

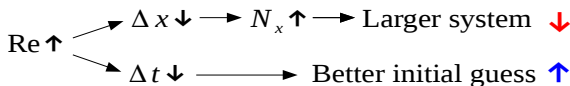
- Will the **complexity** of numerically solving **Poisson's equation** **increase** or **decrease** for **very large scale DNS/LES** simulations of incompressible turbulent flows?

 $Ra = 10^8$  $208 \times 208 \times 400$ **17.5M** $Ra = 10^{10}$  $768 \times 768 \times 1024$ **607M** $Ra = 10^{11}$  $1662 \times 1662 \times 2048$ **5600M**

⁴F.Dabbagh, F.X.Trias, A.Gorobets, A.Oliva. *Flow topology dynamics in a 3D phase space for turbulent Rayleigh-Bénard convection*, **Phys.Rev.Fluids**, 5:024603, 2020. 7 / 17

Smaller and smaller, but how much?

Two competing effects: who (if any) will eventually win?



From classical
K41 theory:

$$\frac{1}{N_x^{K41}} = \frac{\Delta x}{L_x} \sim \frac{\eta}{l} \propto Re^{-3/4}$$

$$\frac{u}{U} \propto Re^{-1/4}$$

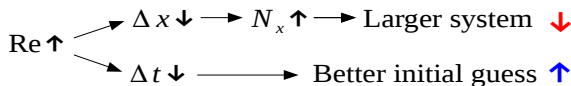
l : biggest eddies (driving scale)
 η : smallest eddies (Kolmogorov length scale)

Question:
 how $\frac{\eta}{l}$ decreases with Re ?

$$\frac{1}{N_t^{K41}} = \frac{\Delta t}{t_{sim}} \sim \frac{t_n}{t_l} \propto \frac{\eta}{l} \frac{U}{u} \propto Re^{-3/4} Re^{1/4} = Re^{-1/2}$$

Smaller and smaller, but how much?

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$$\frac{1}{N_t^{K41}} = \frac{\Delta t}{t_{sim}} \sim \frac{t_\eta}{t_l} \propto \frac{\eta}{l} \frac{U}{u} \propto Re^{-3/4} Re^{1/4} = Re^{-1/2}$$

From CFL condition:

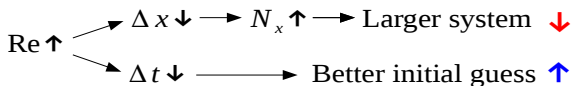
$$\Delta t^{conv} \sim \frac{\Delta x}{U} \quad \Delta t^{diff} \sim \frac{\Delta x^2}{\nu}$$

$$\frac{1}{N_t^{conv}} \sim \frac{\Delta t^{conv}}{t_l} \sim \frac{U}{l} \frac{l}{U} Re^{-3/4} = Re^{-3/4}$$

$$\frac{1}{N_t^{diff}} \sim \frac{\Delta t^{diff}}{t_l} \sim \frac{U}{l} \frac{l^2}{\nu} (Re^{-3/4})^2 = Re^{-1/2}$$

Smaller and smaller, but how much?

Two competing effects: who (if any) will eventually win?



In summary:

$$\frac{1}{N_x^{K41}} = \frac{\Delta x}{L_x} \sim \frac{\eta}{l} \propto \text{Re}^{-3/4}$$

$\alpha = -1/2$ (K41 or diffusion dominated)

$$\frac{\Delta t}{t_l} \sim \text{Re}^\alpha$$

$\alpha = -3/4$ (convection dominated)

Residual of Poisson's equation

$$\nabla^2 p^{n+1} = \frac{1}{\Delta t} \nabla \cdot \vec{u}^p$$

Initial guess $\rightarrow p^n$

$$r^o = \nabla^2 p^n - \frac{1}{\Delta t} \nabla \cdot u^{p,n+1}$$

Residual of Poisson's equation

$$\nabla^2 p^{n+1} = \frac{1}{\Delta t} \nabla \cdot \vec{u}^p$$

Initial guess $\rightarrow p^n$

$$r^o = \nabla^2 p^n - \frac{1}{\Delta t} \nabla \cdot \vec{u}^{p,n+1} = \frac{1}{\Delta t} \nabla \cdot \vec{u}^{p,n} - \frac{1}{\Delta t} \nabla \cdot \vec{u}^{p,n+1} \approx \frac{\partial \nabla \cdot \vec{u}^p}{\partial t} = \nabla \cdot \frac{\partial \vec{u}^p}{\partial t}$$

Residual of Poisson's equation

$$\nabla^2 p^{n+1} = \frac{1}{\Delta t} \nabla \cdot \tilde{\mathbf{u}}^p$$

Initial guess $\rightarrow p^n$

$$r^o = \nabla^2 p^n - \frac{1}{\Delta t} \nabla \cdot \mathbf{u}^{p,n+1} = \frac{1}{\Delta t} \nabla \cdot \mathbf{u}^{p,n} - \frac{1}{\Delta t} \nabla \cdot \mathbf{u}^{p,n+1} \approx \frac{\partial \nabla \cdot \mathbf{u}^p}{\partial t} = \nabla \cdot \frac{\partial \mathbf{u}^p}{\partial t}$$

$$\tilde{r}^o = \nabla^2 \tilde{p}^n - \nabla \cdot \mathbf{u}^{p,n+1} \approx \nabla \cdot \mathbf{u}^{p,n} - \nabla \cdot \mathbf{u}^{p,n+1} \approx \Delta t \frac{\partial \nabla \cdot \mathbf{u}^p}{\partial t} = \Delta t \nabla \cdot \frac{\partial \mathbf{u}^p}{\partial t}$$

Initial guess $\rightarrow \tilde{p}^n = \Delta t p^n$

$$\nabla^2 \tilde{p}^{n+1} = \nabla \cdot \tilde{\mathbf{u}}^p$$

Residual of Poisson's equation

$$\nabla^2 p^{n+1} = \frac{1}{\Delta t} \nabla \cdot \tilde{u}^p$$

Initial guess $\rightarrow p^n$

$$r^o \approx \frac{\partial \nabla \cdot u^p}{\partial t}$$

$$\tilde{r}^o \approx \Delta t \frac{\partial \nabla \cdot u^p}{\partial t}$$

Initial guess $\rightarrow \tilde{p}^n = \Delta t p^n$

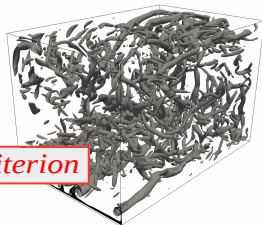
$$\nabla^2 \tilde{p}^{n+1} = \nabla \cdot \tilde{u}^p$$

Residual of Poisson's equation

$$\nabla^2 p^{n+1} = \frac{1}{\Delta t} \nabla \cdot \tilde{u}^p$$

Initial guess $\rightarrow p^n$

Q_G - criterion



$$r^o \approx \frac{\partial \nabla \cdot u^p}{\partial t}$$

What is $\nabla \cdot u^p$?

$$\nabla \cdot u^p = \cancel{\nabla \cdot u^n} - \Delta t \nabla \cdot (u^n \cdot \nabla u^n) + \cancel{v \Delta t \nabla \cdot \nabla^2 u^n} = 2 \Delta t Q_G$$

$$Q_G = -\frac{1}{2} \text{tr}(G^2) \text{ where } G = \nabla u^n$$

$$\tilde{r}^o \approx \Delta t \frac{\partial \nabla \cdot u^p}{\partial t}$$

Initial guess $\rightarrow \tilde{p}^n = \Delta t p^n$

$$\nabla^2 \tilde{p}^{n+1} = \nabla \cdot \tilde{u}^p$$

Residual of Poisson's equation

$$\nabla^2 p^{n+1} = \frac{1}{\Delta t} \nabla \cdot \vec{u}^p$$

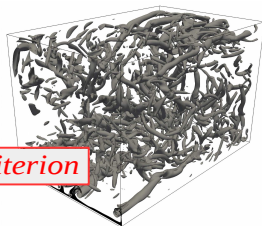
Initial guess $\rightarrow p^n$

$$r^o \approx \frac{\partial \nabla \cdot \vec{u}^p}{\partial t} = 2 \Delta t \frac{\partial Q_G}{\partial t}$$

$$R_G = \det(G) = \frac{1}{3} \text{tr}(G^3)$$

$$\tilde{r}^o \approx \Delta t \frac{\partial \nabla \cdot \vec{u}^p}{\partial t} = 2 \Delta t^2 \frac{\partial Q_G}{\partial t}$$

$$Q_G = -\frac{1}{2} \text{tr}(G^2) \quad \text{where } G = \nabla u^n$$



Exact equations for restricted Euler :

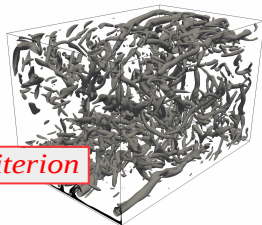
$$\frac{dQ_G}{dt} = -3R_G \longrightarrow \frac{\partial Q_G}{\partial t} = -(u \cdot \nabla) Q_G - 3R_G$$

Residual of Poisson's equation

$$\nabla^2 p^{n+1} = \frac{1}{\Delta t} \nabla \cdot \tilde{u}^p$$

Initial guess $\rightarrow p^n$

Q_G - criterion



$$r^o \approx \frac{\partial \nabla \cdot u^p}{\partial t} = 2 \Delta t \frac{\partial Q_G}{\partial t} \approx -2 \Delta t \{ (u \cdot \nabla) Q_G + 3 R_G \}$$

$$\tilde{r}^o \approx \Delta t \frac{\partial \nabla \cdot u^p}{\partial t} = 2 \Delta t^2 \frac{\partial Q_G}{\partial t} \approx -2 \Delta t^2 \{ (u \cdot \nabla) Q_G + 3 R_G \}$$

Exact equations for restricted Euler :

$$\frac{dQ_G}{dt} = -3 R_G \longrightarrow \frac{\partial Q_G}{\partial t} = -(u \cdot \nabla) Q_G - 3 R_G$$

Residual of Poisson's equation

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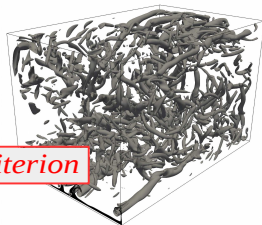
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$$r^o \approx \frac{\partial \nabla \cdot u^p}{\partial t} = 2 \Delta t \frac{\partial Q_G}{\partial t} \approx -2 \Delta t \{ (u \cdot \nabla) Q_G + 3 R_G \}$$

$$\tilde{r}^o \approx \Delta t \frac{\partial \nabla \cdot u^p}{\partial t} = 2 \Delta t^2 \frac{\partial Q_G}{\partial t} \approx -2 \Delta t^2 \{ (u \cdot \nabla) Q_G + 3 R_G \}$$

Initial guess $\rightarrow \tilde{p}^n = \Delta t p^n$

$$\nabla^2 \tilde{p}^{n+1} = \nabla \cdot \tilde{u}^p$$



Q_G - criterion

Residual of Poisson's equation in Fourier space

In summary:

$$r^o \approx \frac{\partial \nabla \cdot u^p}{\partial t} = 2 \Delta t^p \frac{\partial Q_G}{\partial t} \approx -2 \Delta t^p \{ (u \cdot \nabla) Q_G + 3 R_G \}$$

$$\nabla^2 \tilde{p}^{n+1} = \frac{1}{\Delta t} \nabla \cdot \tilde{u}^p$$

$$p = \{1, 2\}$$

$$\nabla^2 \tilde{p}^{n+1} = \nabla \cdot \tilde{u}^p$$

$$\frac{\Delta t}{t_l} \sim \text{Re}^\alpha \begin{cases} \alpha = -1/2 & (\text{K41 or diffusion dominated}) \\ \alpha = -3/4 & (\text{convection dominated}) \end{cases}$$

$$\frac{1}{N_x^{\text{K41}}} = \frac{\Delta x}{L_x} \sim \frac{\eta}{l} \propto \text{Re}^{-3/4}$$

Residual of Poisson's equation in Fourier space

In summary:

$$r^o \approx \frac{\partial \nabla \cdot u^p}{\partial t} = 2 \Delta t^p \frac{\partial Q_G}{\partial t} \approx -2 \Delta t^p \{ (u \cdot \nabla) Q_G + 3 R_G \}$$

$$\nabla^2 p^{n+1} = \frac{1}{\Delta t} \nabla \cdot \tilde{u}^p$$

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$$\nabla^2 \tilde{p}^{n+1} = \nabla \cdot \tilde{u}^p$$

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(inertial range)

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Residual of Poisson's equation in Fourier space

In summary:

$$r^o \approx \frac{\partial \nabla \cdot u^p}{\partial t} = 2 \Delta t^p \frac{\partial Q_G}{\partial t} \approx -2 \Delta t^p \{ (u \cdot \nabla) Q_G + 3 R_G \}$$

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$$\|r^n\|^2 = \int_1^{k_{max}} (\hat{\omega}_k^n \hat{r}_k^0)^2 dk \approx \int_1^{Re^{3/4}} \hat{\omega}_k^{2n} Re^{2\tilde{\alpha}} k^{2\beta} dk$$

$$\hat{\omega} = \frac{\hat{r}_k^{n+1}}{\hat{r}_k^n}$$

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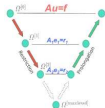
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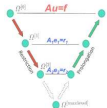
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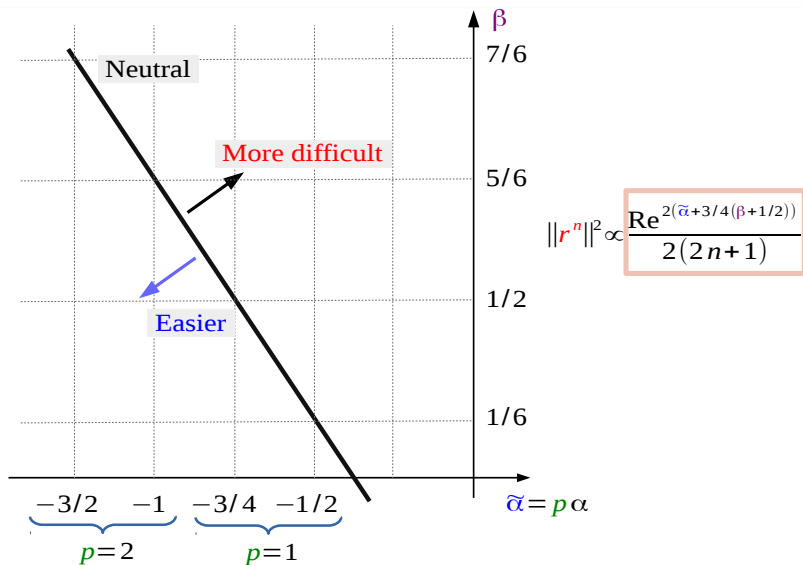
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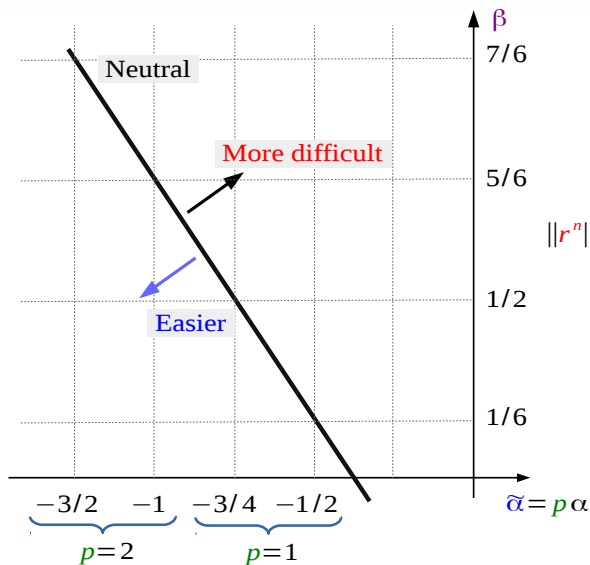
Solver convergence

$\{\tilde{\alpha}, \beta\}$ phase space



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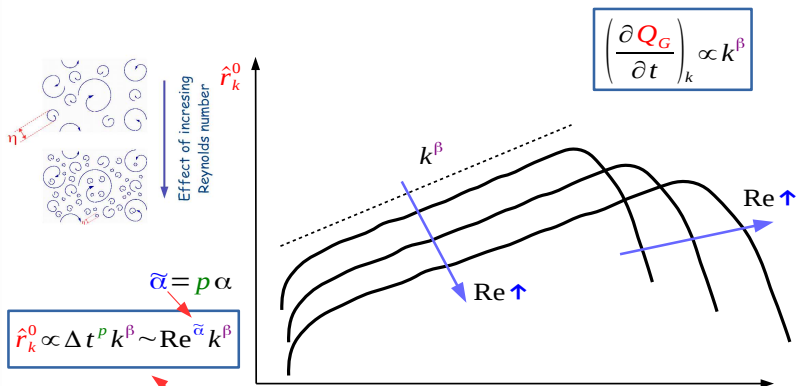


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Solver convergence

$\{\tilde{\alpha}, \beta\}$ phase space



Two competing effects!!!

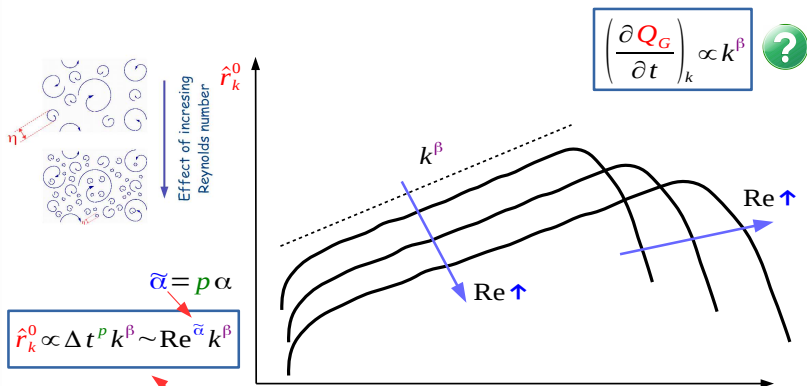
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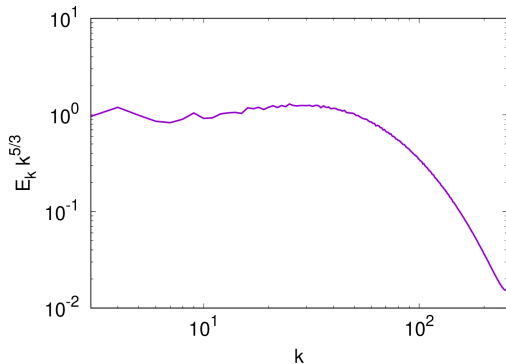
Homogeneous isotropic turbulence

Kolmogorov theory predictions

$$E_k = C_k \varepsilon^{2/3} k^{-5/3}$$

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Homogeneous isotropic turbulence

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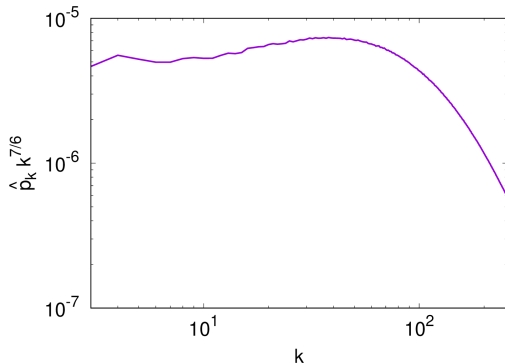
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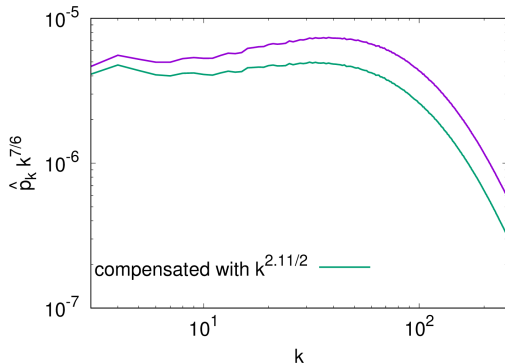
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Homogeneous isotropic turbulence

New derivations

$$E_k = C_k \varepsilon^{2/3} k^{-5/3}$$

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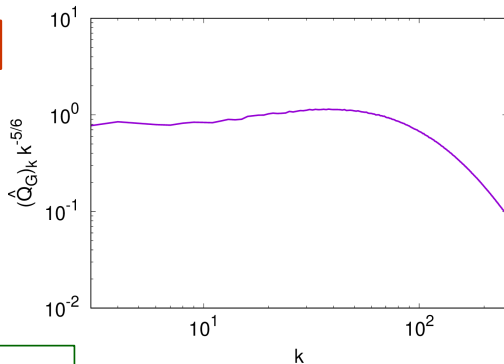
$$\hat{p}_k \propto k^{-7/6}$$

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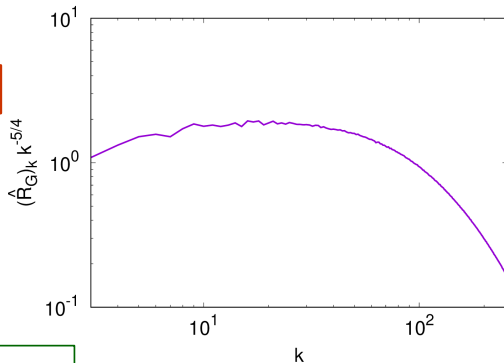
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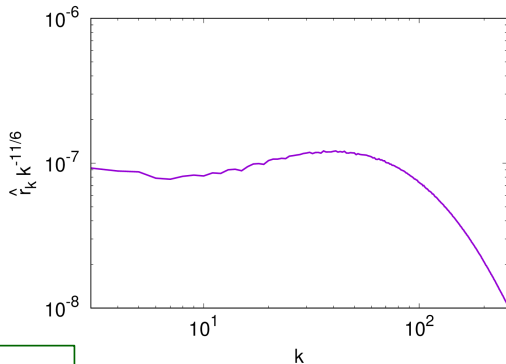
$$(\hat{r})_k \propto k^{5/6+1} = k^{11/6}$$

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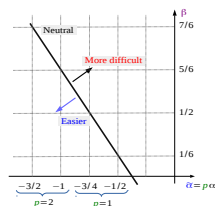


Concluding remarks

- **Two competing effects** on the convergence of Poisson's equation have been identified.

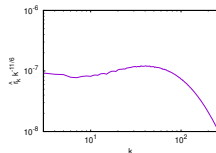
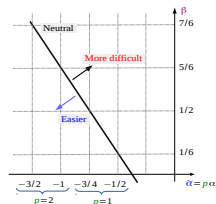
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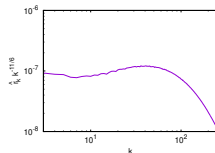
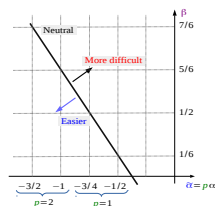
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On-going and near future research:

- Carrying out simulations at higher Re_λ
- Extending the analysis to more complex flows

Thank you for your attendance